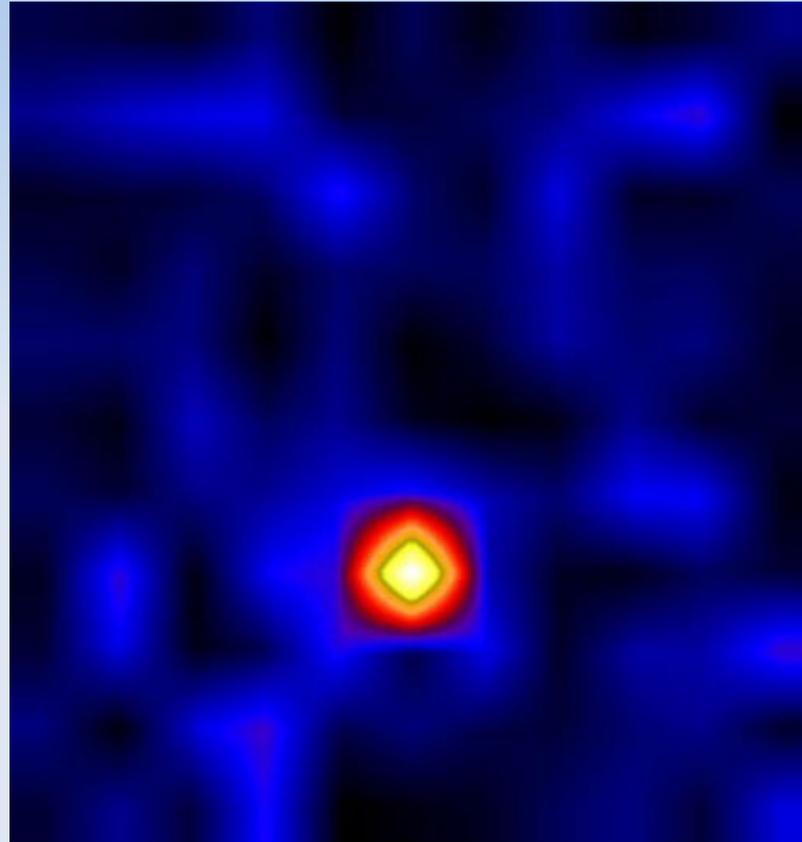


Compact Stars



Lecture 6

Summary of the previous lecture

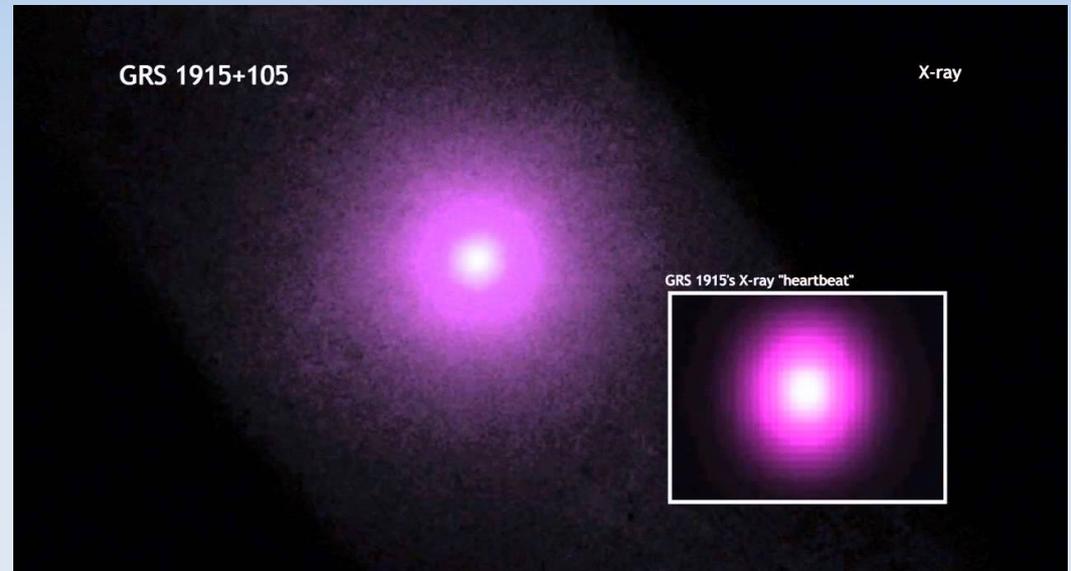
- I presented the systems with accreting white dwarfs – cataclysmic variables - and their classification.
- I showed that about the time-dependent modeling of accretion disk thermal-viscous instabilities, induced by the partial hydrogen ionisation can explain the dwarf novae outbursts. The opacity change in disk affects thermal balance and leads to instabilities.
- The dominant radiation pressure (in black hole or neutron star X-ray binaries) may lead to similar cycles, for hotter disks. In the radiation pressure dominated disk some stabilisation is needed (to form an upper branch on the stability "S-curve"). It is provided by the advective model (slim disk).
- The instabilities of the accretion flows manifest in quasi-periodic oscillations of luminosity. Examples: dwarf novae (best fitted with two values of viscosity parameter), microquasars.

Microquasars, quasars, and other jet sources

- Today we will start from microquasars
- Some of them exhibit the radiation pressure instability in action, and oscillations in the accretion disks at high accretion rates (close to Eddington limit)
- Their primary feature are the relativistic blobs of plasma ejected along the symmetry axis
- They are similar to quasars, and other astrophysical jet sources, but on stellar mass scale.

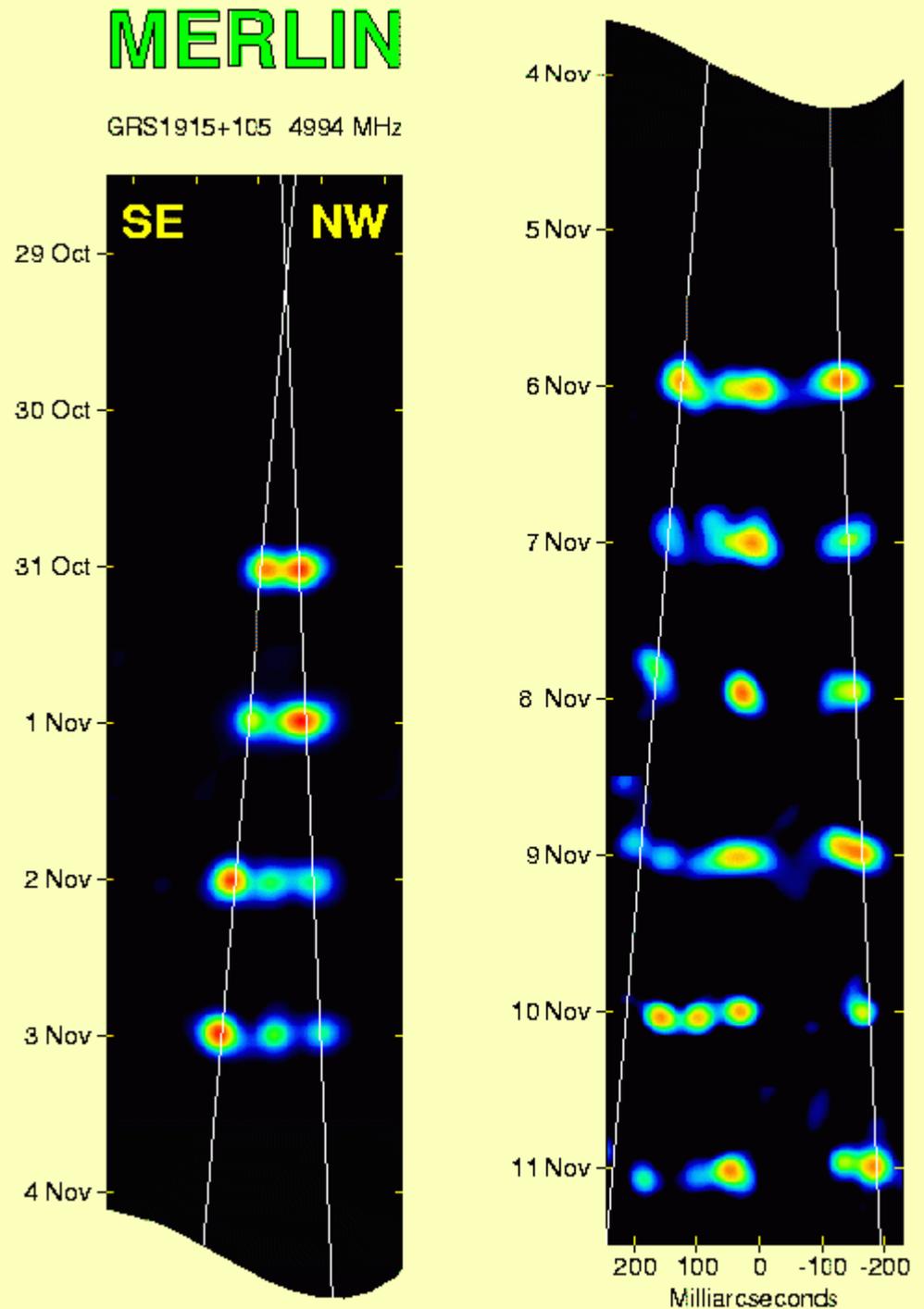
GRS 1915+105

- Accreting black hole, almost Eddington luminosity
- Radiation pressure instability in action: limit-cycle oscillations of X-ray luminosity
- Exhibits radio jets: analogous to quasars
- Interesting source to study jet-disk connection



GRS 1915+105

Radio observations
from 1994
(revealed apparent
superluminal motion)



Apparent superluminal motion

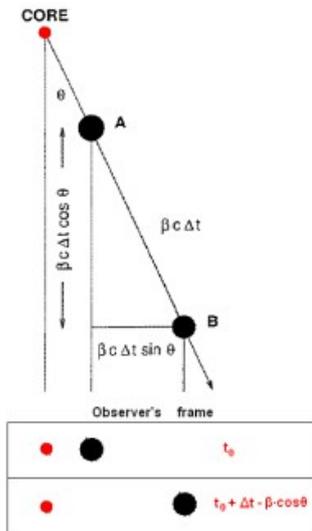


Figure 3. Cartoon illustrating the superluminal motion phenomenon. From R.A. Laing [14].

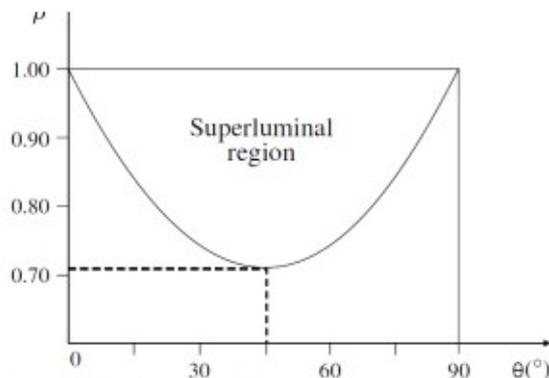


Figure 4. Superluminal region for a linearly expanding source, shown as an area on the $\beta - \theta$ diagram. From D.F. Falla & M.J. Floyd [15].

- $\Delta t_{\text{obs}} = \Delta t_{\text{em}} - d/c$

- $d = v \Delta t_{\text{em}} \cos \theta$

$$v_{\text{obs}} = v \sin \theta \Delta t_{\text{em}} / \Delta t_{\text{obs}}$$

$$= \beta c \sin \theta / (1 - \beta \cos \theta)$$

$$\beta_{\text{obs}} = 1.26 \text{ for } \beta = 0.92 \text{ and } \theta = 70^\circ$$

Microquasars

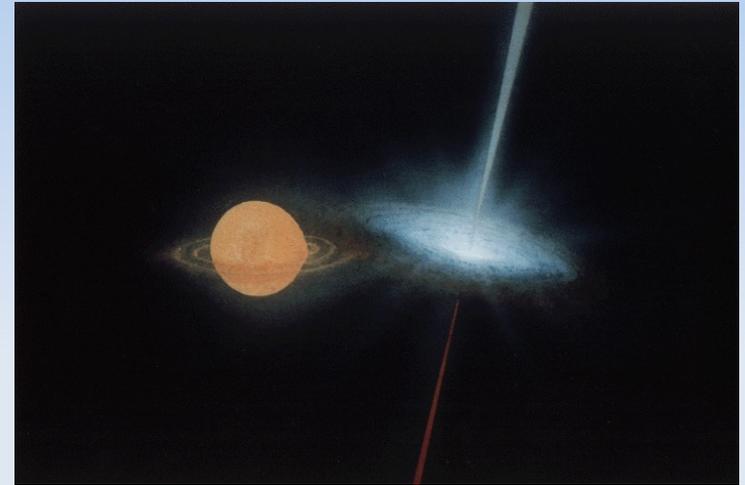
Known from 1990-ties

Smaller “cousins” of quasars

Characteristics:

- Strong and variable radio emission
 - Presence of jets
- Large X-ray luminosity

The energetics is explained by a black hole accretion



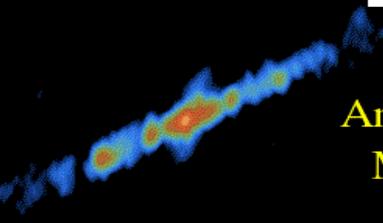
Simultaneous radio-X-ray observations allowed the identification of radio jets with X-ray binaries (GRANAT and VLBA)

Many are known: over 30 sources including Integral data (as of 2017)

(<http://www.aim.univ-paris7.fr/CHATY/Microquasars/microquasars.html>)

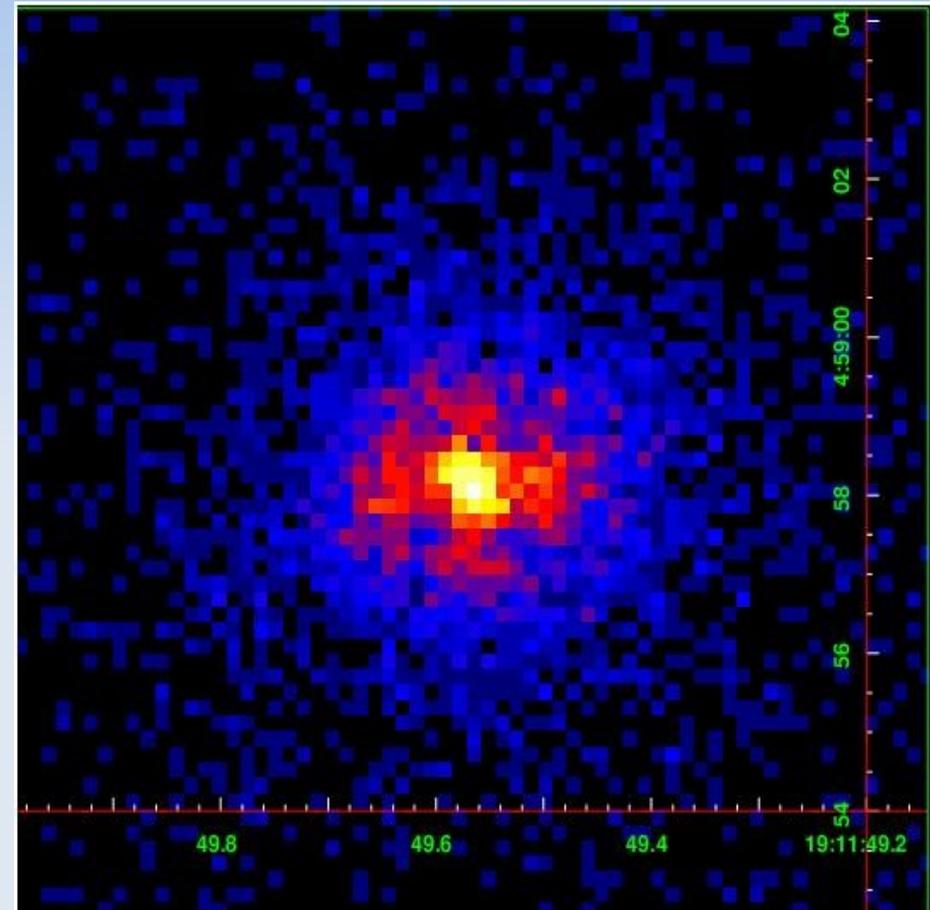
Microquasars: prototype of the class

SS433
VLBA



Astronomy logos: NRAO, Earth, Associated Universities, Inc., NSF

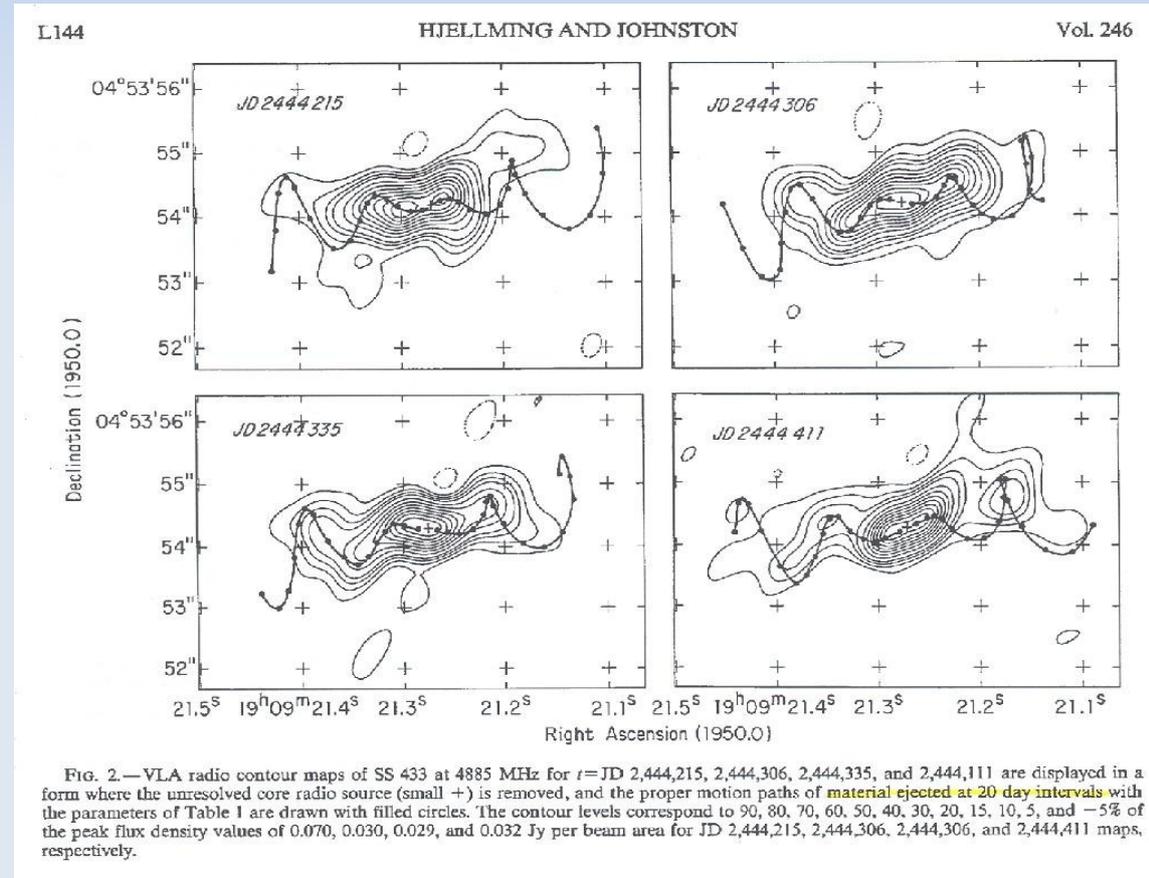
Amy Mioduszewski
Michael Rupen
Craig Walker
Greg Taylor



SS 433 – first microquasar discovered (Margon et al. 1979). Chandra X-ray image

SS 433 in radio

- Microquasar with NS
- Detected by VLA (Hjellming & Johnston 1981)
- Mildly relativistic, $v \sim 0.26 c$
- Precessing jet model proposed



Quasars: definition



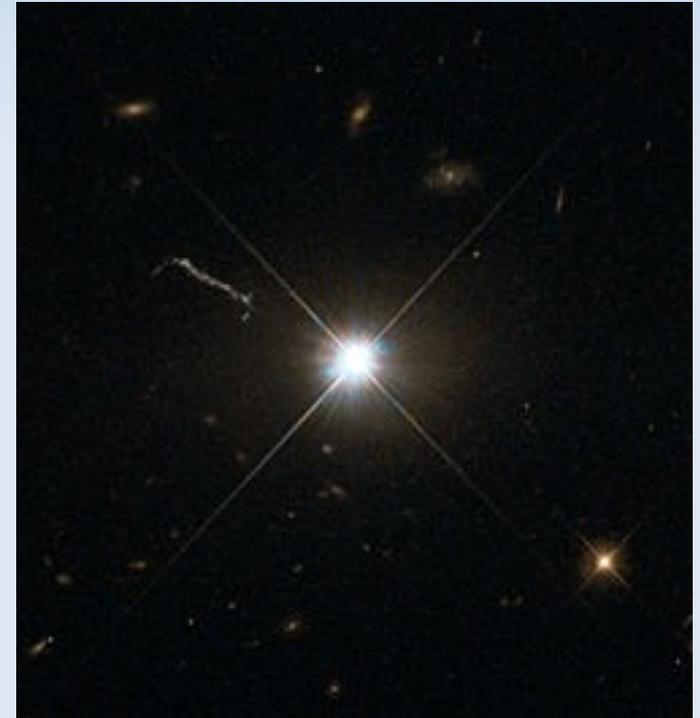
Allan Sandage

Discovered in 1960-ties
Quasi-Stellar Radiosources
Currently about 1 000 000
known (many discovered
SDSS survey)

About 10% of them are radio
loud

Matthews T., Sandage A., et al.
(1963, ApJ): Optical identification
of 3C 48

Schmidt M. & Oke B.
(Nature, 1963):
Z=0.158 in 3C 273



3C 273 – image from
Hubble Telescope

Energy source of quasars

Quasars shine at all electromagnetic spectrum, from IR to X. (and Radio i Gamma)

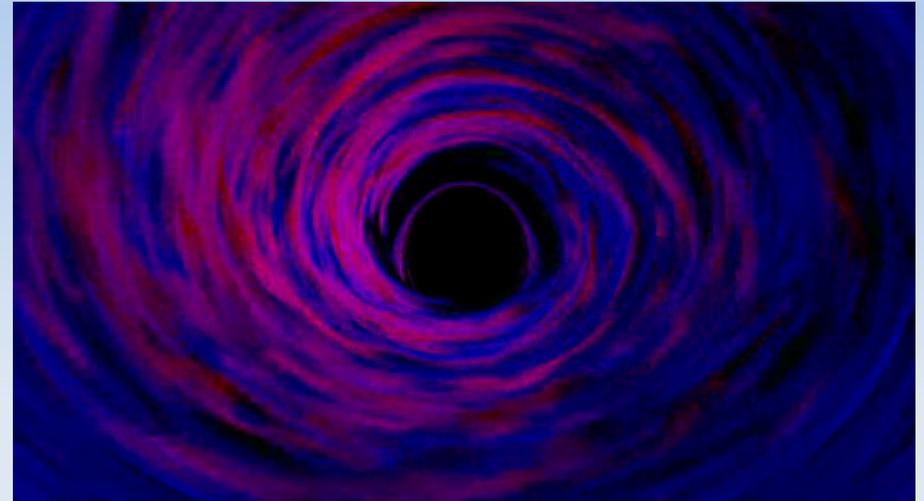
The luminosity of a QSO can be 100x higher than a galaxy

Fast variability of flux indicates that the emitting region is very compact (spatial scale of a Solar System)

Size and timescale proportional to M_{BH}

$$R_{\text{schw}} = 2GM_{\text{BH}}/c^2 ;$$

$$\Delta T \sim M_{\text{BH}}$$



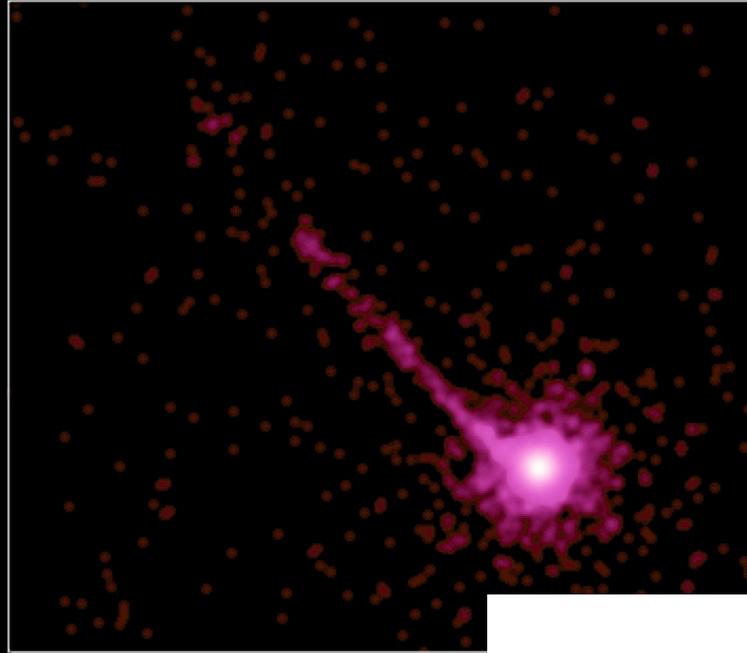
Accretion onto black hole
 $E_p = GMm/R$

Efficiency $E_p/mc^2 = GM/Rc^2$

Quasars emit jets

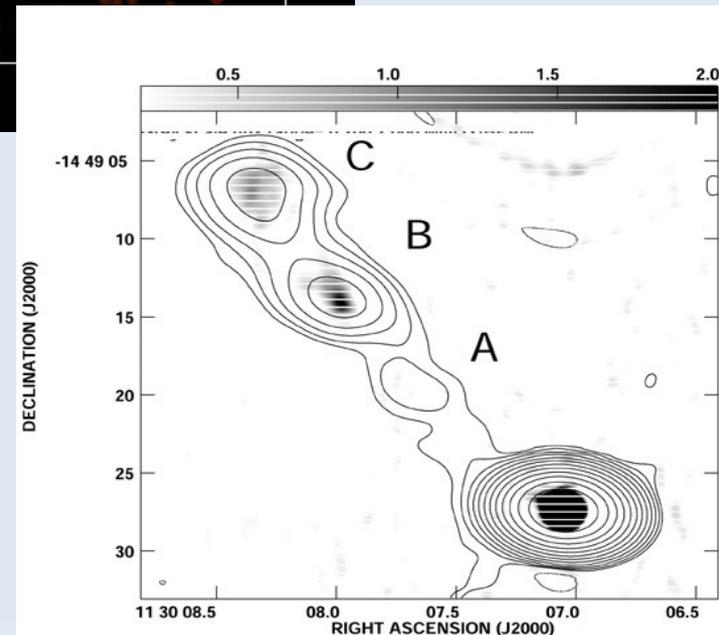


PKS 1127-145



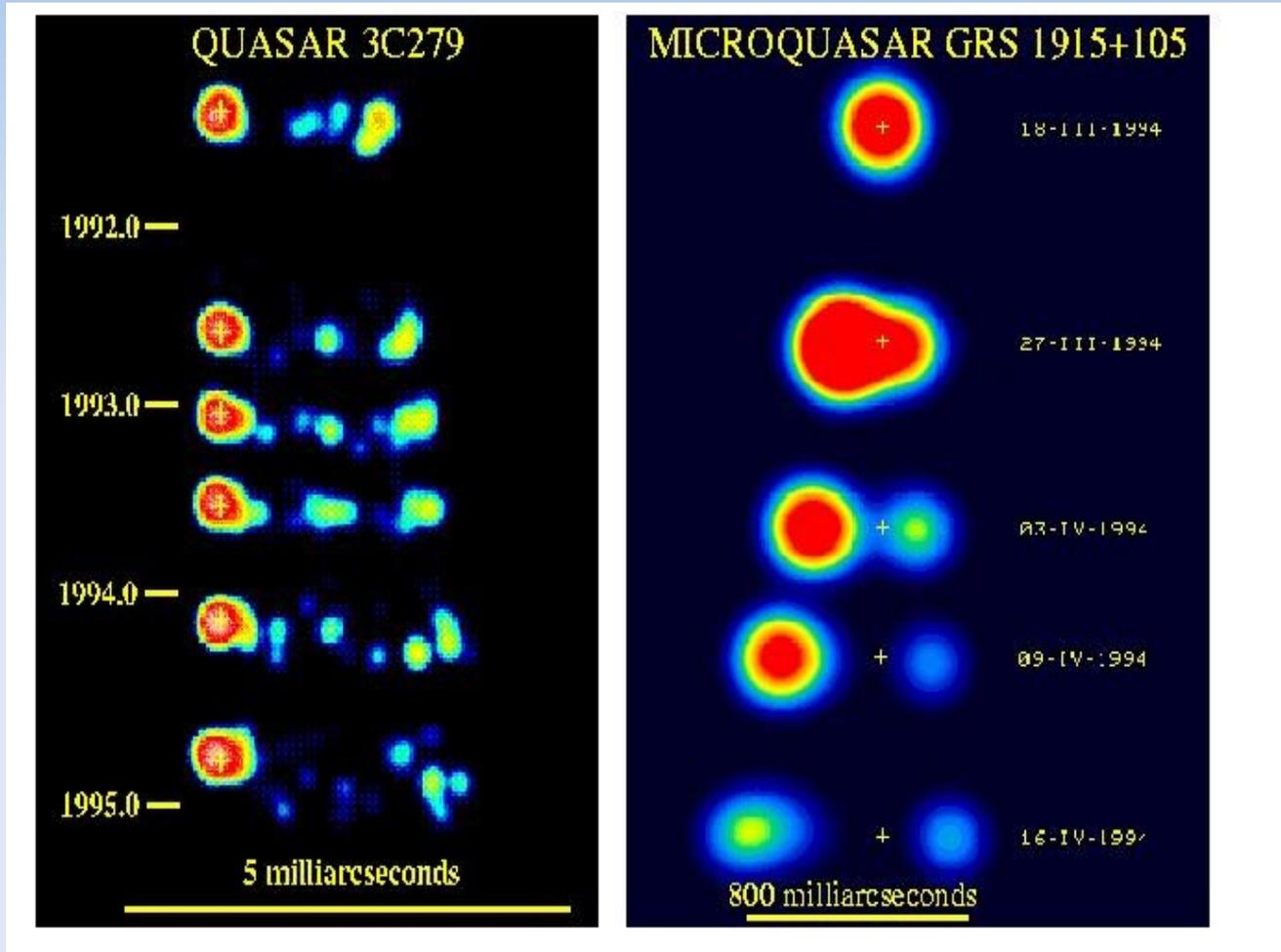
**Quasar PKS 1127-145,
X-ray image from Chandra
(*Siemiginowska & Bechtold, 2000*)
Optical (HST), radio (VLA)**

**Large scale jet has size of million
light years**



Microquasars

- Defined by the analogy to quasars
- Radio observations of jets show apparently superluminal motions

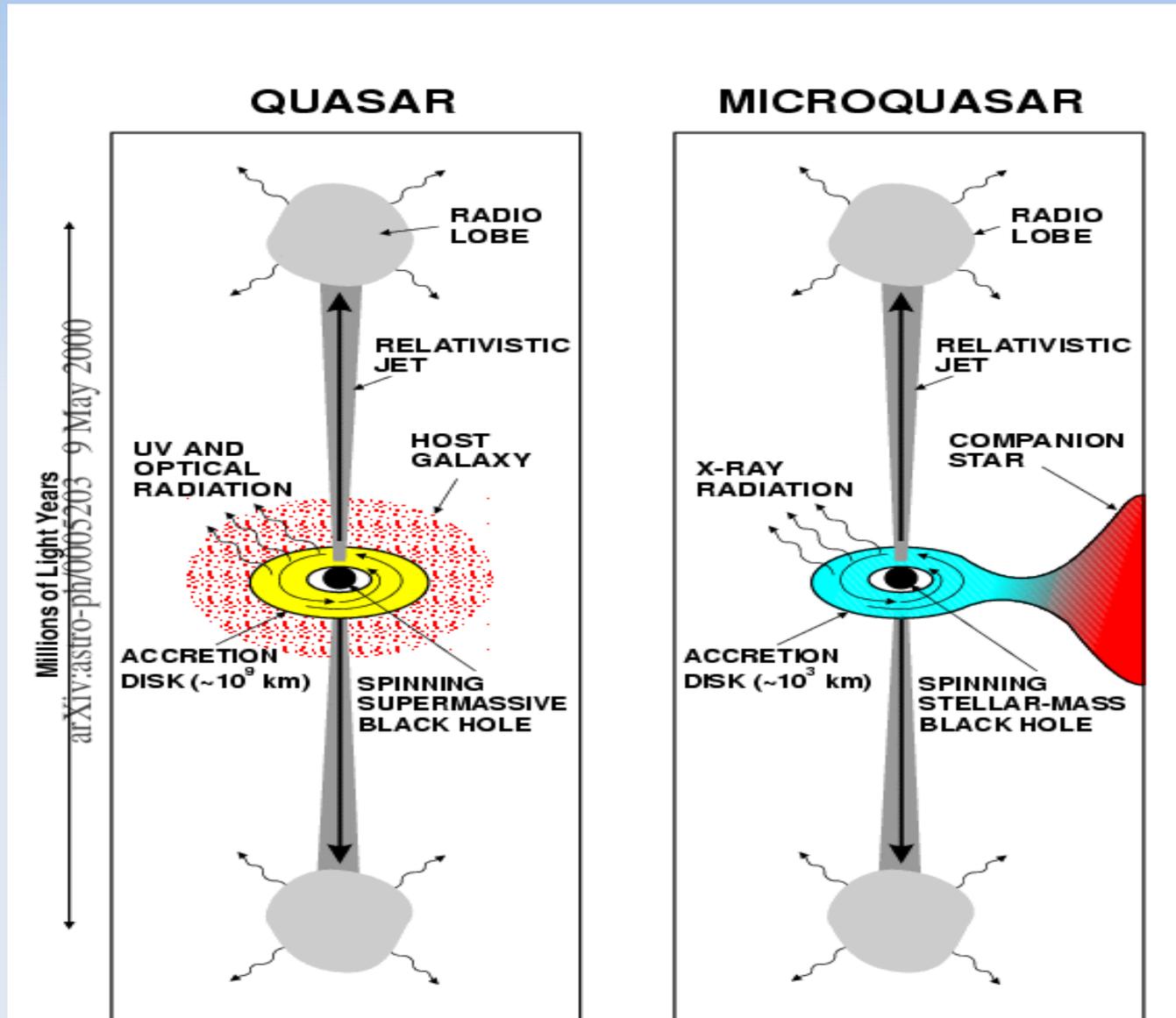


Mirabel & Rodriguez, 1994, Nature

Comparison QSOs vs. μ QSOs

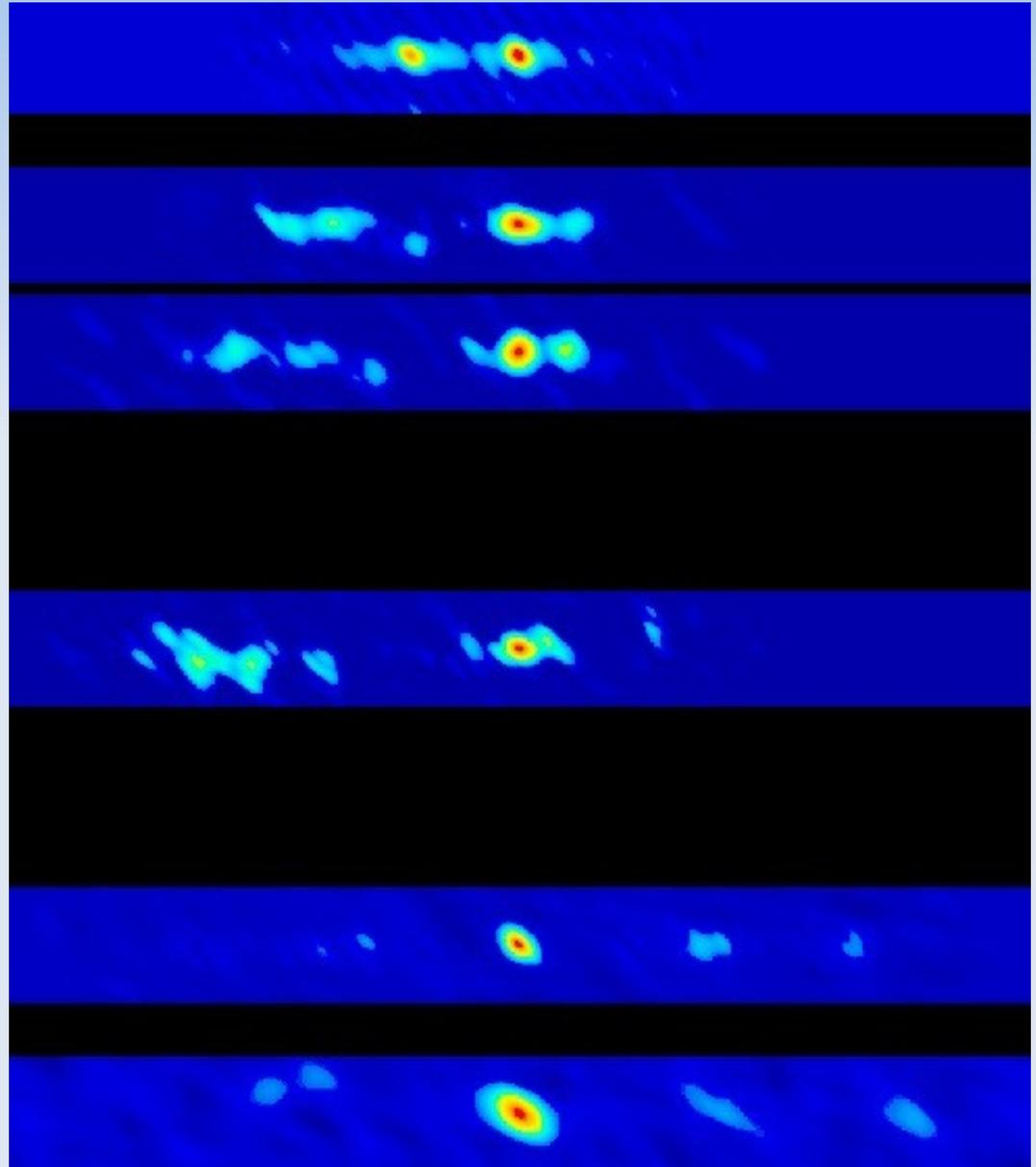
- Point radiosources, discovered in 1960's
- Active nuclei of galaxies
- Mass of BH $10^8 - 10^9 M_{\text{Sun}}$
- $T_{\text{disk}} \sim 10^5 \text{ K}$
- Luminosity $\sim 10^{47} \text{ erg/s}$
- Accretion rate $\sim 10 M_{\text{sun/yr}}$
- Variability in timescale of years
- Jet range about 10^7 light years
- Discovered in our Galaxy in 1990's
- Binary stars systems
- Mass of BH/NS $2 - 20 M_{\text{Sun}}$
- $T_{\text{disk}} \sim 10^7 \text{ K}$
- Luminosity $\sim 10^{37} \text{ erg/s}$
- Accretion rate $\sim 10^{-9} M_{\text{sun/yr}}$
- Variability in timescales of minutes
- Jet range of several light years

Comparison of Quasars and Microquasars



Second superluminal BH XRB

- GRO 1655-40
 - microquasar
 - VLBA radio data
 - Apparently superluminal motions
 - Observed after the X-ray outburst



Expanding blobs

- From observations we know proper motions of the approaching and receding blobs (in arc seconds per day). This gives the expansion velocity times the cosine of the angle between jet and our line of sight.
- Knowing the distance to the source, we can independently determine the jet velocity and viewing angle.
- The jet velocity can be also measured from the observations of emission lines.

Synchrotron radiation

- High brightness temperatures, ‘nonthermal’ spectra and polarisation measurements indicate an origin as synchrotron emission from relativistic electrons

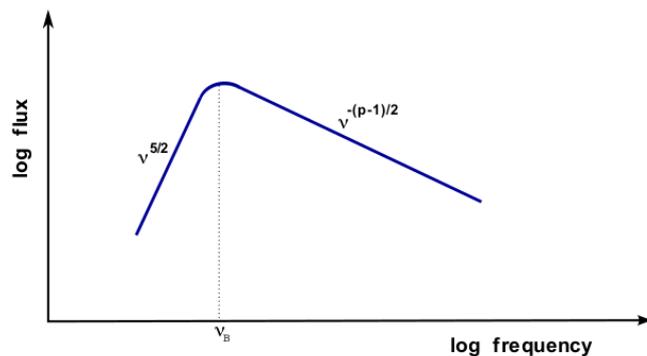
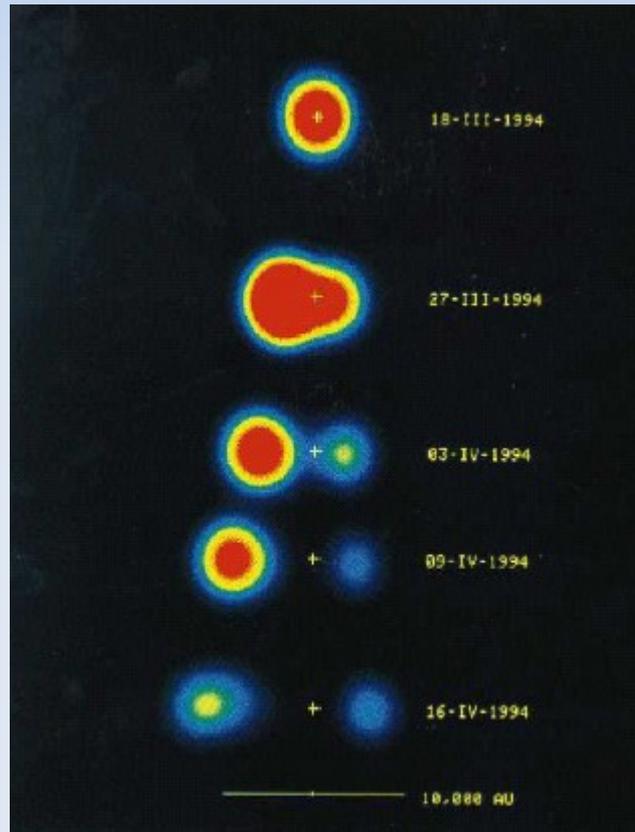


Figure 2: Total synchrotron spectrum - Shape for electron power law $n(\gamma) \propto \gamma^{-p}$, with $j_\nu \propto \nu^{5/2}$ in the low-frequency, optically-thick regime ($\nu < \nu_B$), and $j_\nu \propto \nu^{-(p-1)/2}$ in the optically thin ($\nu > \nu_B$) regime. Note that for an electron distribution with low- and high-energy cut-off γ_{min} , γ_{max} , the optically thin spectrum could be more complex with low-frequency part $j_\nu \propto \nu^{1/3}$, mid-frequency part $j_\nu \propto \nu^{-(p-1)/2}$ and high-energy frequency part $j_\nu \propto e^{-\nu/\nu_{max}}$, $\nu_{max} \simeq \gamma_{max}^2 \nu_L$.

Observed optically thin spectral indices $-0.4 \geq \alpha \geq -0.8$, indicate $1.8 \leq p \leq 2.6$ for the underlying electron distribution

Doppler boosting



$$S_{r,a} = S_0 / [\gamma (1 \pm \beta \cos\theta)]^{k+\alpha}$$

$S_0 \propto \nu^{-\alpha}$ radiation flux density in the source frame;

$k=2$ for steady jet or $k=3$ for 'blobby jet"

For small angles and large luminosities, the apparent brightness of approaching blob is boosted, (a), and for receding blob it is weakened, (r)

Inverse Compton process

Total radiated IC power for single electron (Thomson regime, $\epsilon_i = \gamma h\nu \ll m_e c^2$)

$$P_{IC} = \frac{4}{3} \gamma^2 \beta^2 \sigma_T c U_{rad}$$

This is similar to single particle synchrotron power:

$$P_{syn} = \frac{4}{3} \gamma^2 \beta^2 \sigma_T c U_B$$

Reason: In both cases, electron is accelerated by an electric field which it observes in its instantaneous rest-frame (electromagnetic wave/photon, $\mathbf{E}' = \mathbf{v} \times \mathbf{B}$). Electron does not care about true origin of electric field.

Spectral index of scattered photon spectrum:

$$P_{IC}(\nu) \propto \nu^{-\alpha}$$

off power-law electron distribution with index p , $n_e(\gamma) \propto \gamma^{-p}$, is

$$\alpha = (p - 1)/2$$

This is similar to synchrotron as well.

Compton upscattering of synchrotron photons by the electrons themselves (SSC) preserves the spectral index in the Thomson regime

$$\alpha_{syn} \approx \alpha_{IC}$$

Two main radiative processes

Ratio of the **Inverse Compton** to the **Synchrotron** flux density, for a fixed magnetic field strength, as a function of the jet viewing angle.

Depends on the jet bulk Lorentz factor and synchrotron power-law spectral index α (1.1 and 0.6, with dashed and solid lines).

See "Radiative processes in astrophysics", Rybicki & Lightman

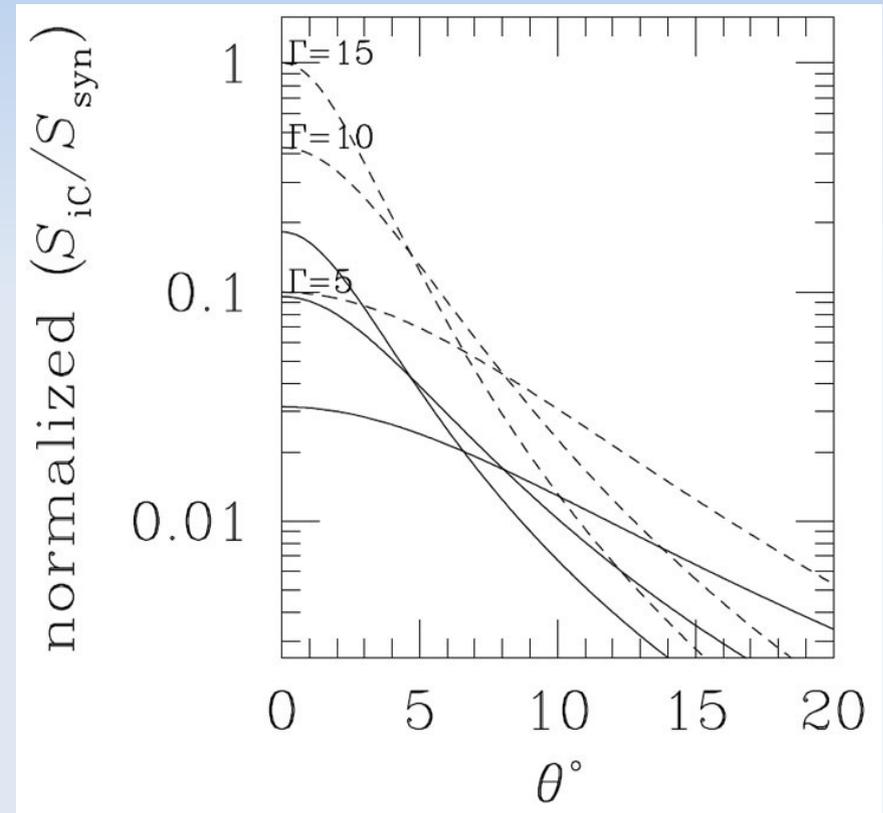


Figure from Worrall & Birkinshaw, 2006, "Lecture Notes in Physics"

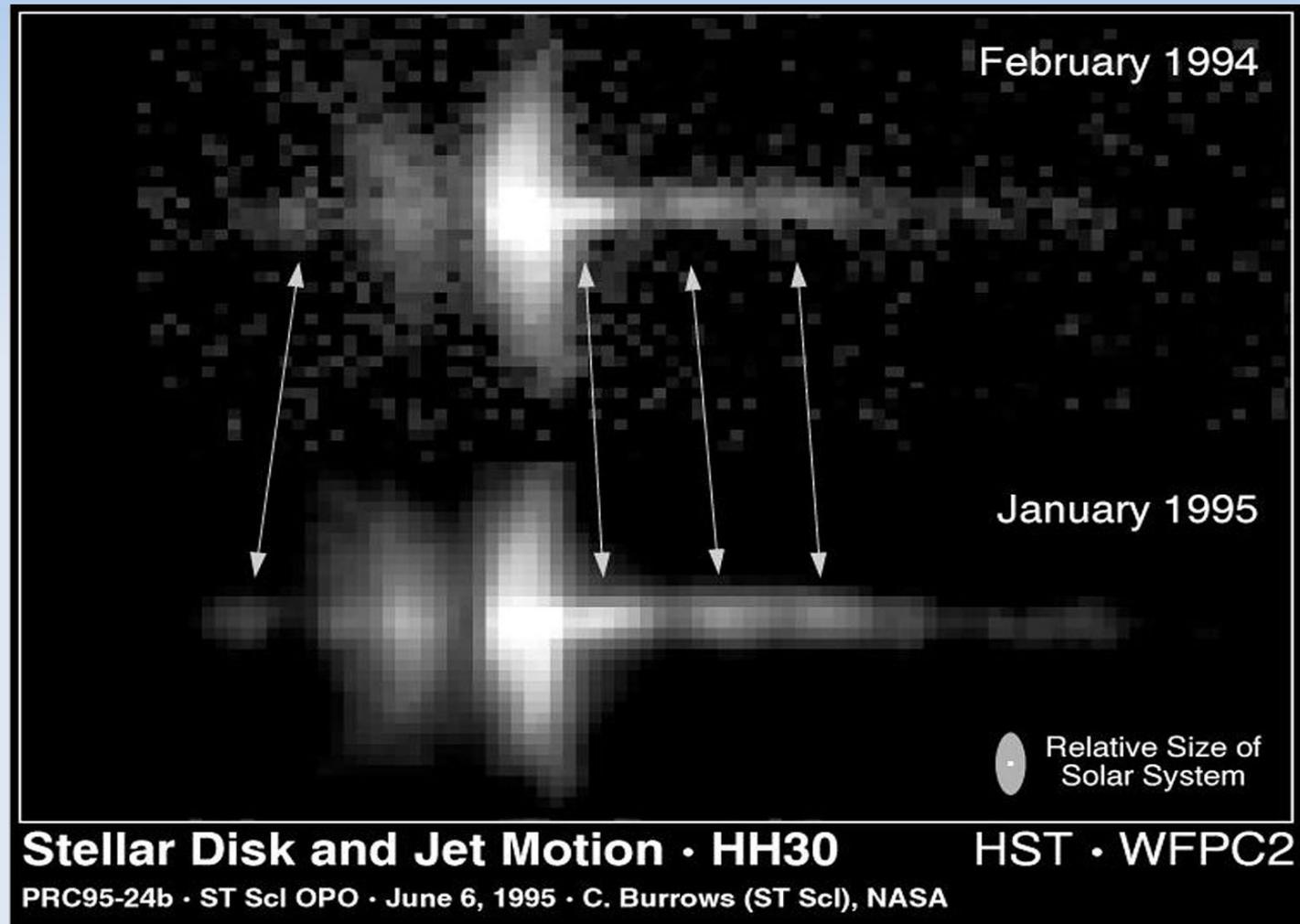
Break

Systems producing jets

- Jets are observed in a number of types of sources:
 - Young stellar objects
 - Massive X-ray binaries
 - Black hole X-ray transients
 - Low mass X-ray binaries
 - Symbiotic stars
 - Planetary nebulae
 - Supersoft X-ray sources
 - Active galactic nuclei

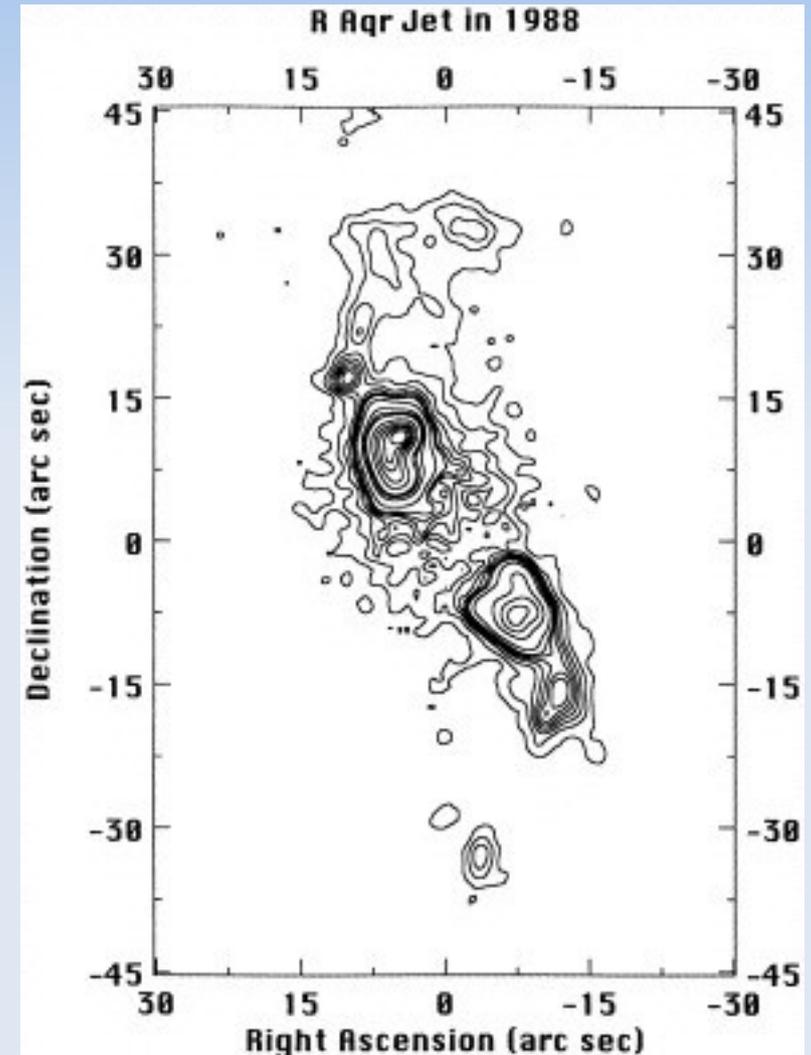
Examples: YSO

- Herbig-Haro object
- Young star surrounded by a thin disk
- $V_{\text{jet}} = 100\text{-}350 \text{ km/s}$



Examples: R Aqr

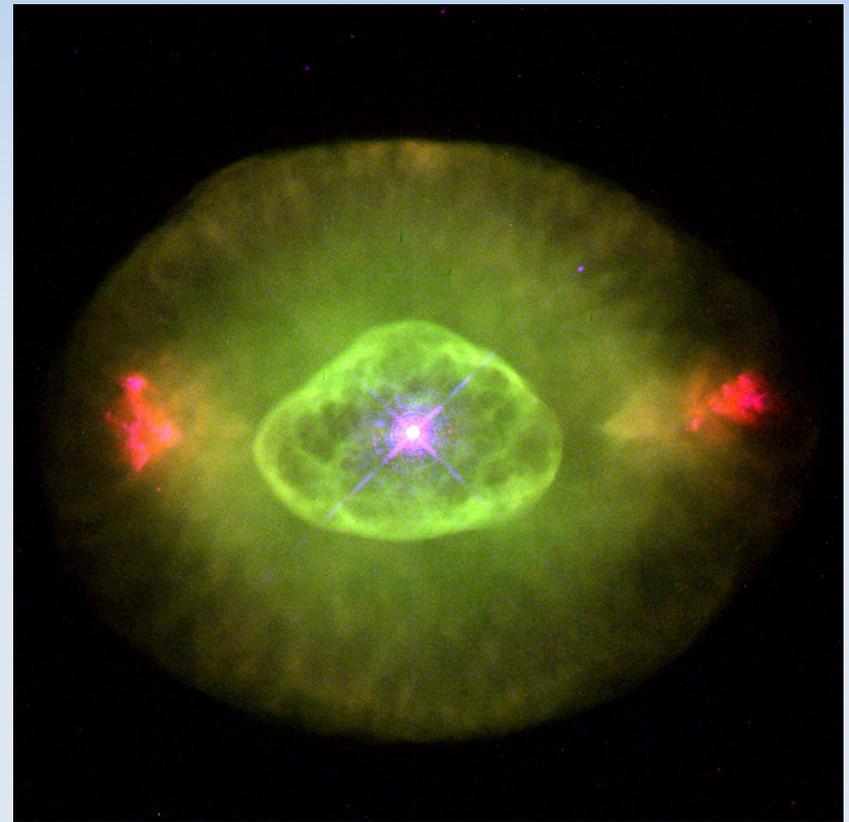
- Symbiotic binary
- Contains white dwarf and Mira star (variable red giant)
- Jet in optical and radio bands (Burgarella et al. 1992)
- Estimated jet age is ~100 years



[O III] 5007 intensity image of the R Aqr symmetrical jet shown as a contour plot (Hollis et al. 1998).

Examples: planetary nebulae

- FILIERS: fast low ionization features are observed in some planetary nebulae (e.g. Redman & Dyson, 1999)
- Interpreted as recombination zones behind supersonic fronts produced in mass loaded jets
- $V \sim 200\text{-}1000$ km/s



NGC 6826

Example: Cir X-1

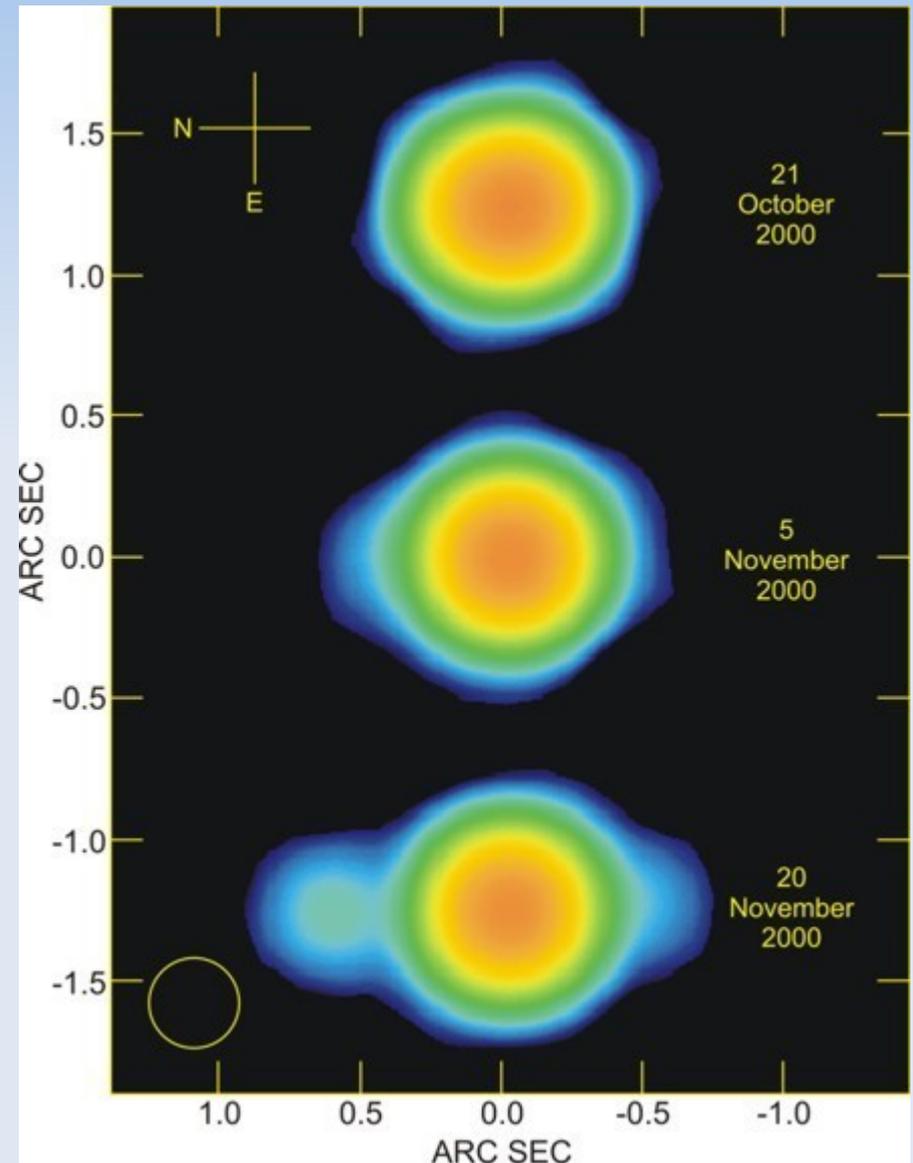
- Low mass X-ray binary with neutron star
- Very young NS (4600 yrs, in supernova remnant; "bouncing baby")
- Jet discovered in 2007
- The jet might be twisted and precessing (Coriat et al. 2019)



Chandra image

Examples: HMXB binaries

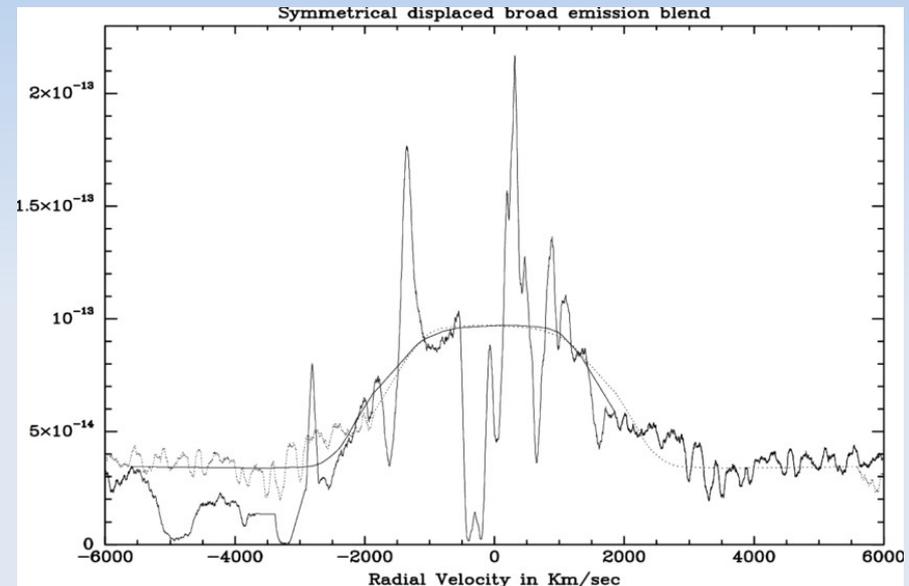
- Cyg X-3
- Contains compact star on $2.4 M_{\text{sun}}$ and Wolf-Rayet star
- Jet emitted at viewing angle about 14 deg (Koljonen et al. 2019)



(NRAO image)

Examples: supersoft X-ray sources

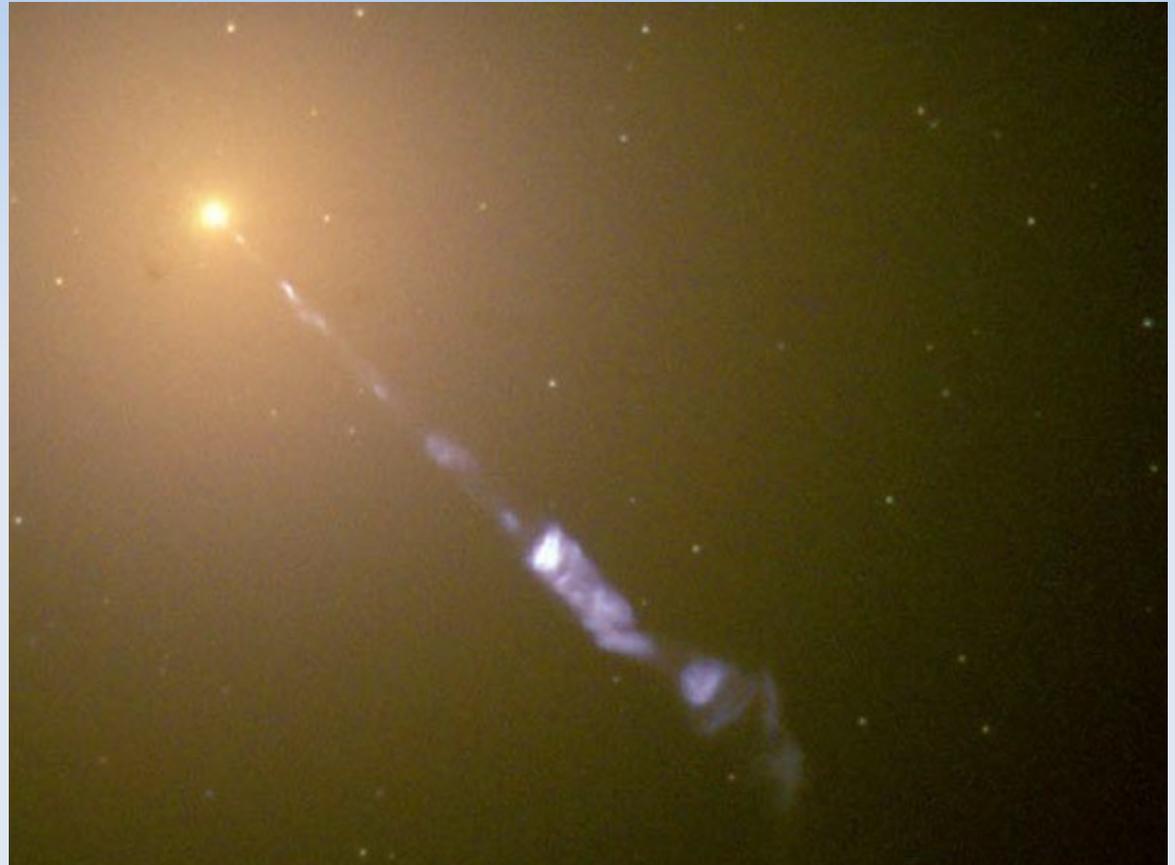
- RX J0513.9-6951
- Compact binary in Large Magellanic Cloud
- Accreting white dwarf; bb temperature 10^5 K, luminosity $L \sim 10^{38}$ erg/s
- Discovered by ROSAT All sky survey
- Unusual emission lines, interpreted as high velocity bipolar outflows
- $V \sim 4000$ km/s (Crampton et al. 1996)



Moving broad absorption feature near 1020A as possible O VI outflow associated with a precessing jet.

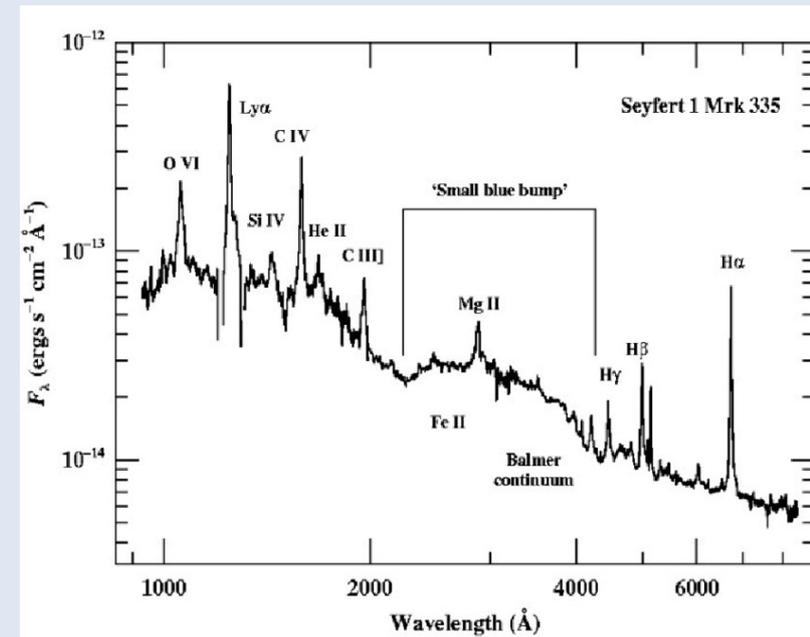
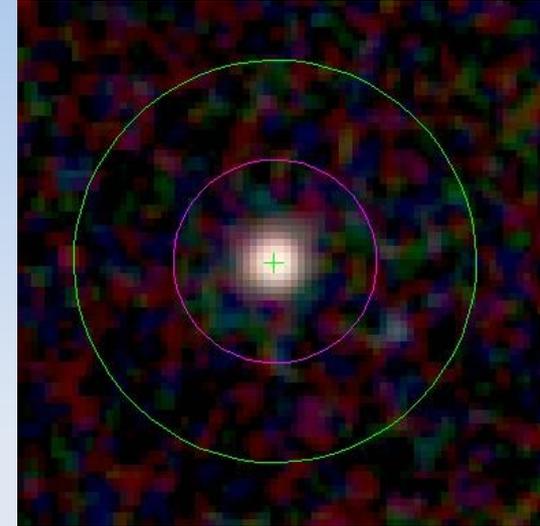
Examples: AGN

- M87
- Supergiant, elliptical galaxy
- Optical jet (HST image)
- Velocity of jet: $\gamma \sim 3$
- Large scale: 5000 light years



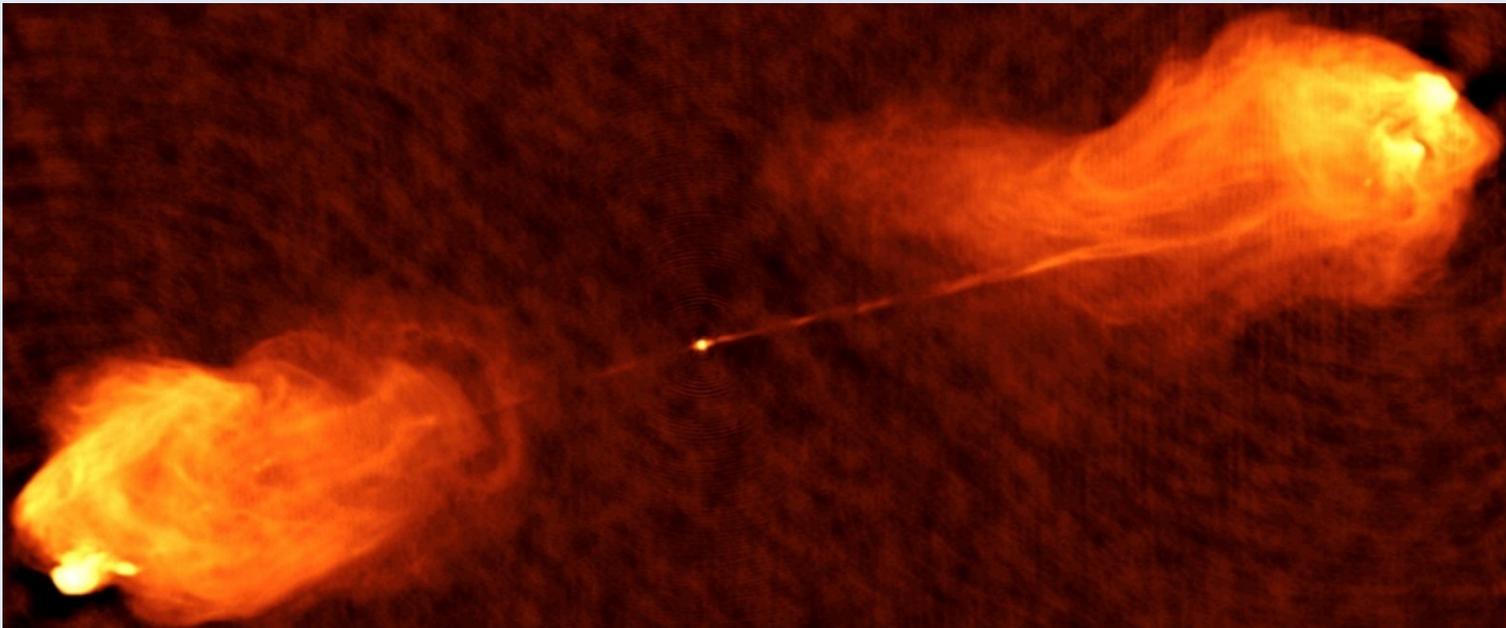
Active galactic nuclei

- An active galactic nucleus (AGN) is a compact region at the center of a galaxy that emits a significant amount of energy across the electromagnetic spectrum.
- The luminosity is not produced by stars. Such excess, non-stellar emissions have been observed in the radio, microwave, infrared, optical, ultra-violet, X-ray and gamma ray wavebands.
- A galaxy hosting an AGN is called an active galaxy.

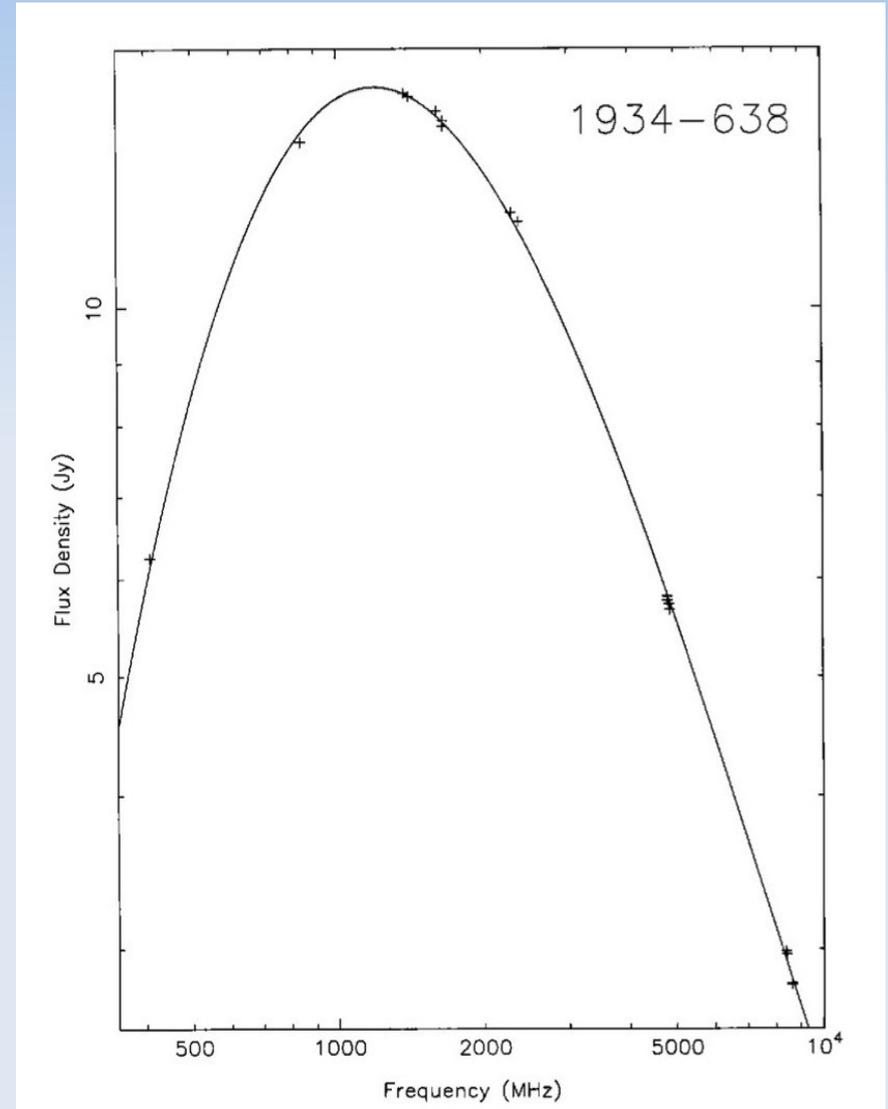
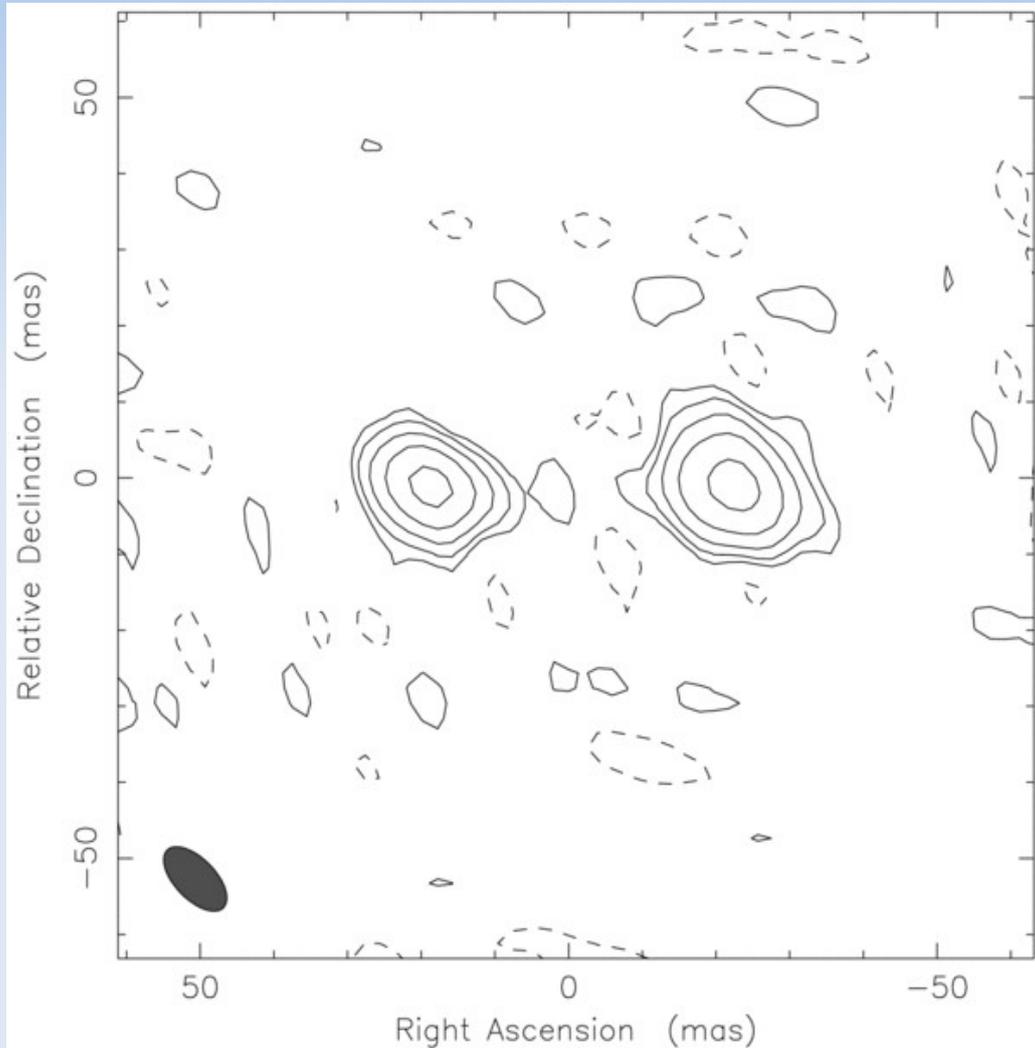


Examples: radiosources

- Jet velocity $\Gamma \sim 10$
- Cygnus A: FR II type radiosource (Fanaroff-Riley), characteristic bright spots and faster motion of jets



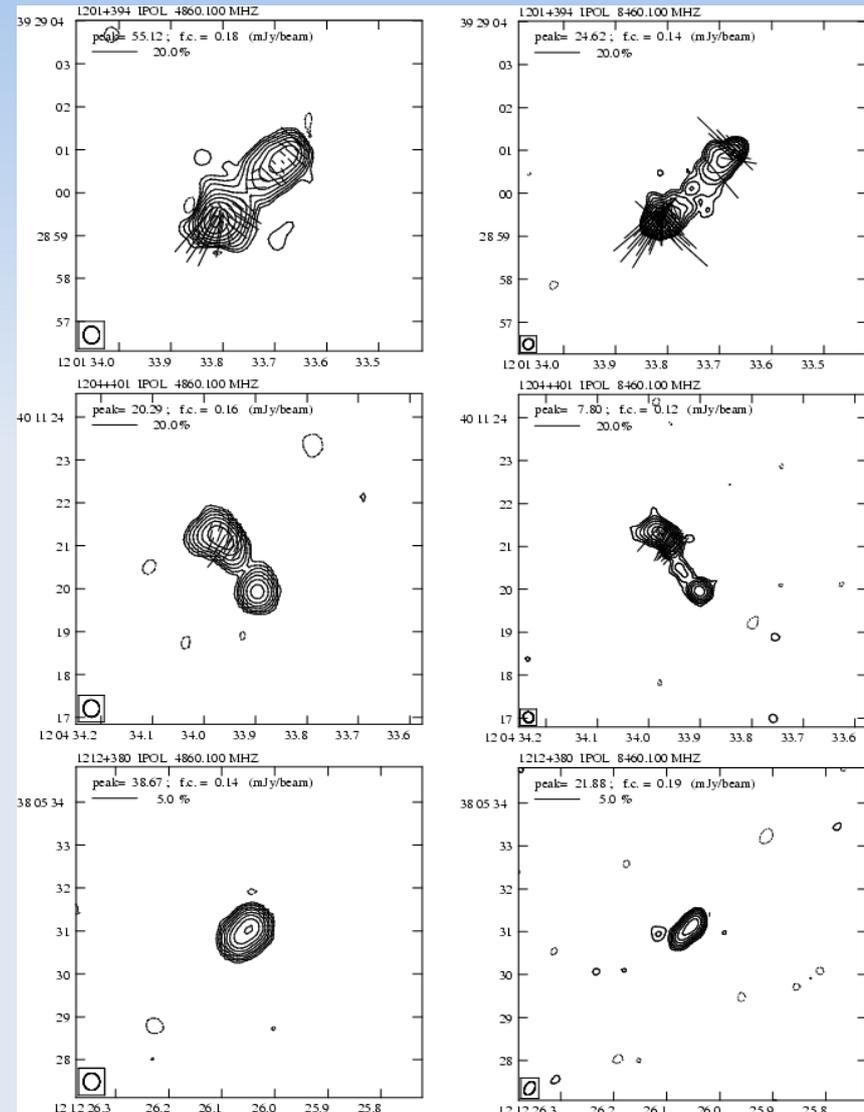
GPS/CSS sources



- PKS 1934-638

GPS radiosources

- Sources appear to be young, because they are compact
- Statistically, their young age would not match with the number of mature radioquasars
- This supports the hypothesis of reactivation

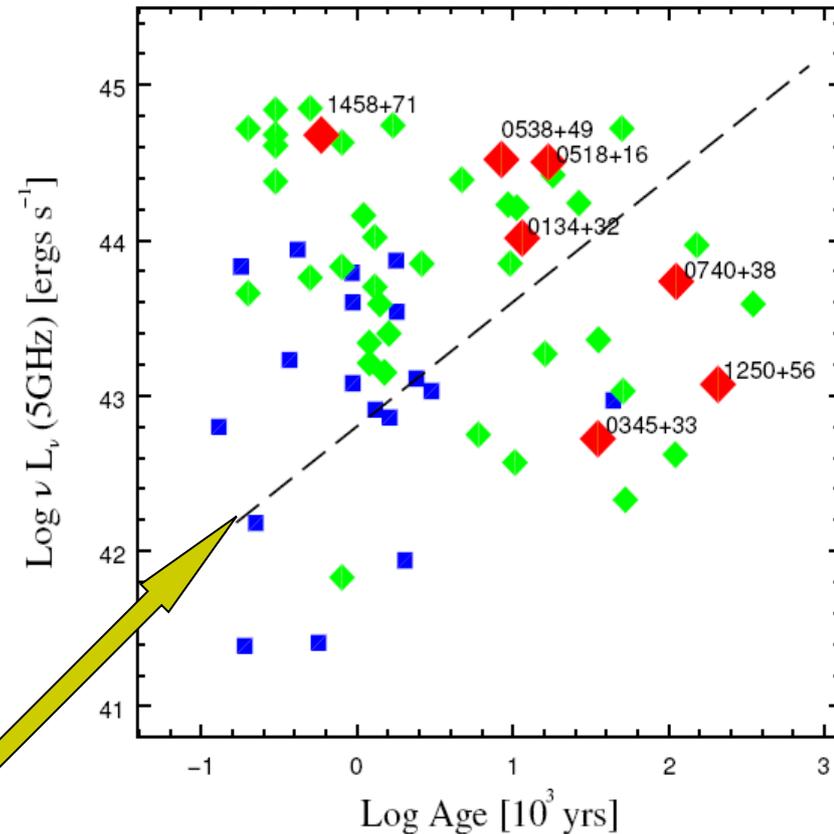


Radio sources: ageing estimates

- The sources are bright and have small size → their age is below 10^4 yrs.
- Observed statistical overabundance of the compact radio sources, in comparison with the large scale > 1 kpc sources (O'Dea & Baum 1997)
- In fact they may be reactivating every 10^4 yrs (Reynolds & Begelman 1997)

GPS radiosources

- Giga-Hertz Peak Spectrum radioquasars
- Sample of 70 apparently young sources
- Hypothesis of reactivation due to the duty-cycle



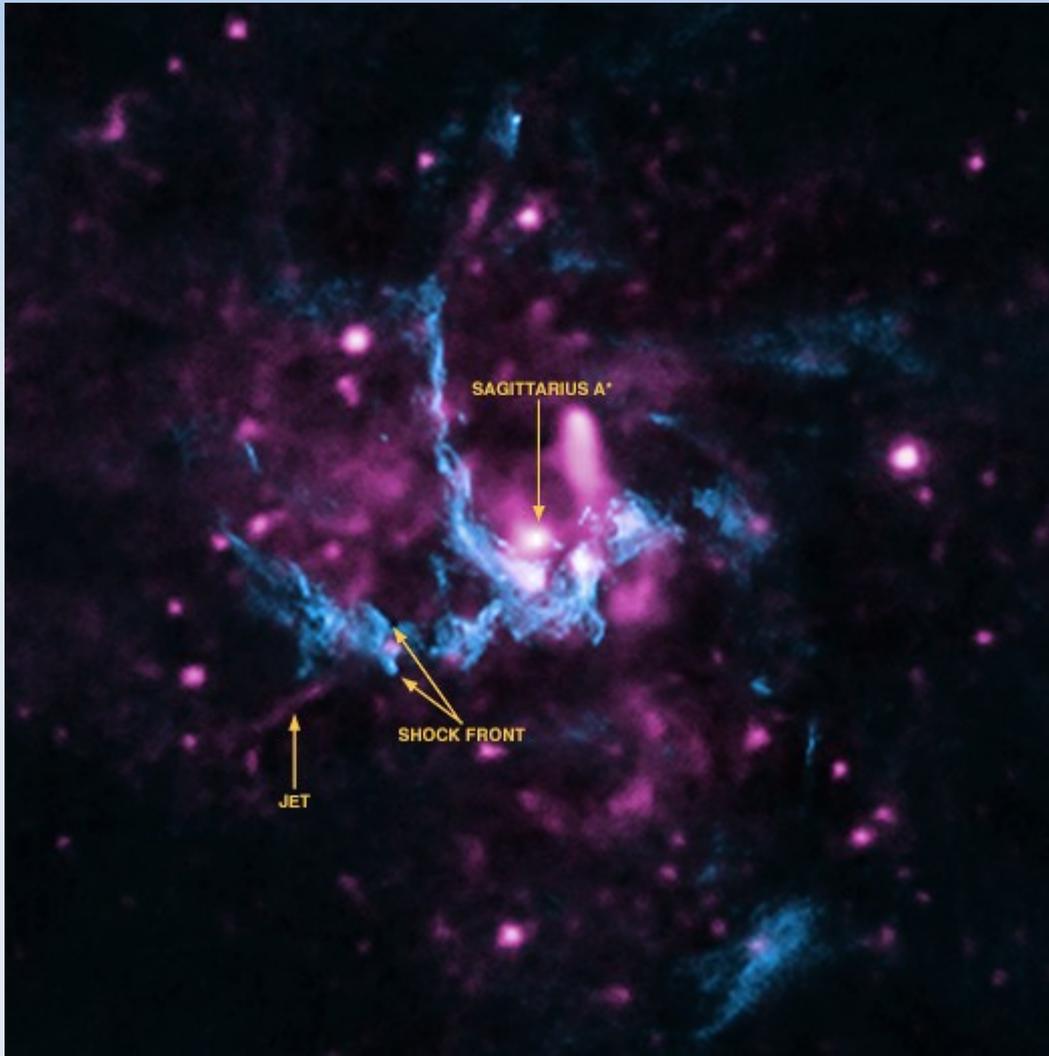
$$\log \left(\frac{T_{\text{burst}}}{\text{yr}} \right) \approx 1.25 \log \left(\frac{\nu L_{\nu} (5\text{GHz})}{\text{erg/s}} \right) + 0.38 \log \left(\frac{\alpha}{0.02} \right) + 1.25 \log K_{5\text{GHz}} - 53.6.$$

Our Galaxy Center



Chandra image of Sgr A*

Does Sgr A* produce a jet?



Chandra X-ray image of the feature, given the name of G359.944-0.052 by its Galactic coordinates, exhibits a needle-shape.

Attributed to jet structure

(Z. Li, Morris, Baganoff, 2013, ApJ)

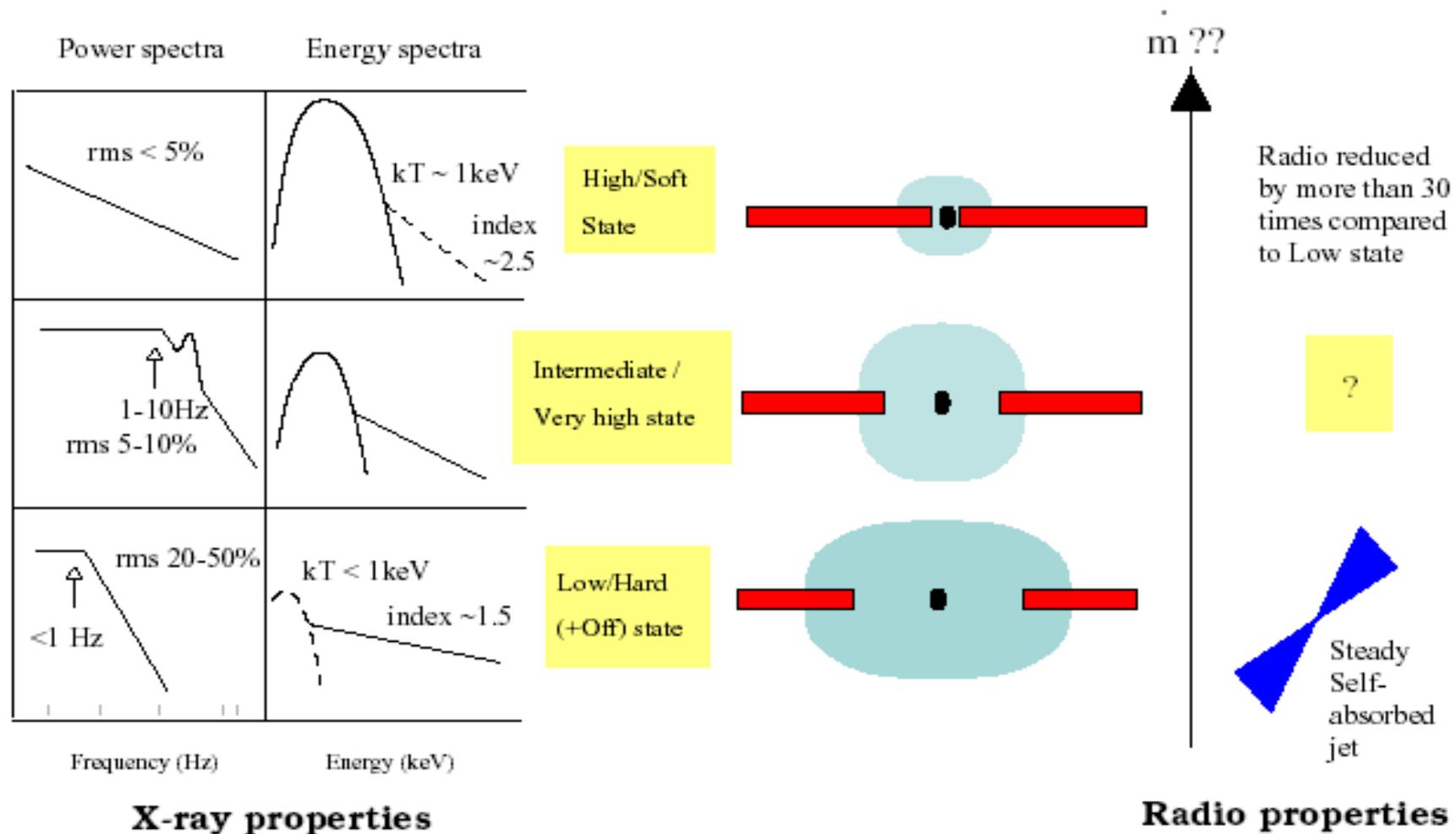
Origin of jets

- In all cases, the estimated jet velocity of order of escape velocity from the central object
- The jet must therefore be produced in the vicinity of the central object, at inner regions of an accretion disk
- In microquasars, observed correspondence between the X-ray, IR and Radio fluxes and constant time delay suggest that the inner disk ejects the plasma, responsible then for IR and Radio flares by synchrotron emission (e.g. Fender et al. 1997)

Advection Dominated Accretion (ADAF)

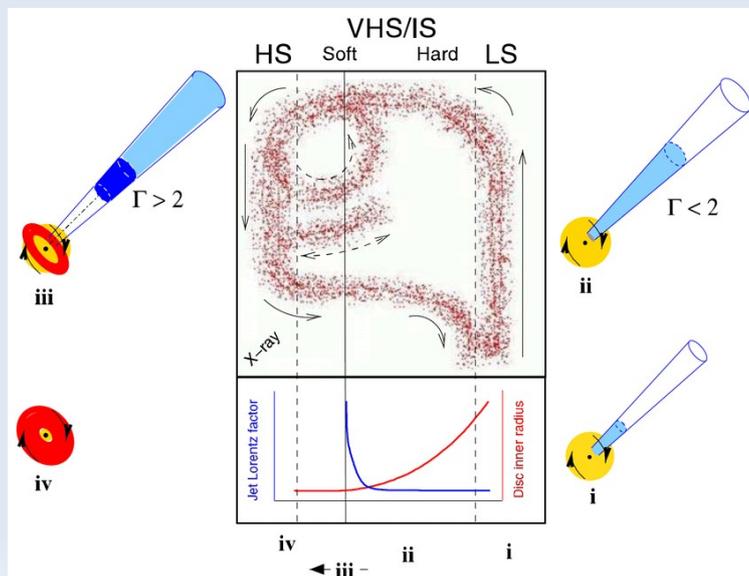
- Infalling (or accreting) matter falls into a black hole without radiating away much of its energy, less than one percent.
- Density of the infalling matter is low, so the particles make few collisions on their way beyond the event horizon.
- If the matter is falling onto a neutron star, or is spiralling slowly in a thick disk toward a black hole, it can radiate away ten to a hundred times more energy than in the advection case.

Phenomenology: X-ray - radio



ADAF and jets

- Radio emission is too bright to be produced by thermal electrons on the accretion flow
- It may come from nonthermal electrons in a jet
- Radio jets are observed in hard states of the XRB
- Optically thin ADAF might be associated with jet production, in low luminosity states



Established pattern:

- LS → steady jet
- HS → no jet
- VHS → transient jet

Fender et al. (2005)

Various accretion solutions

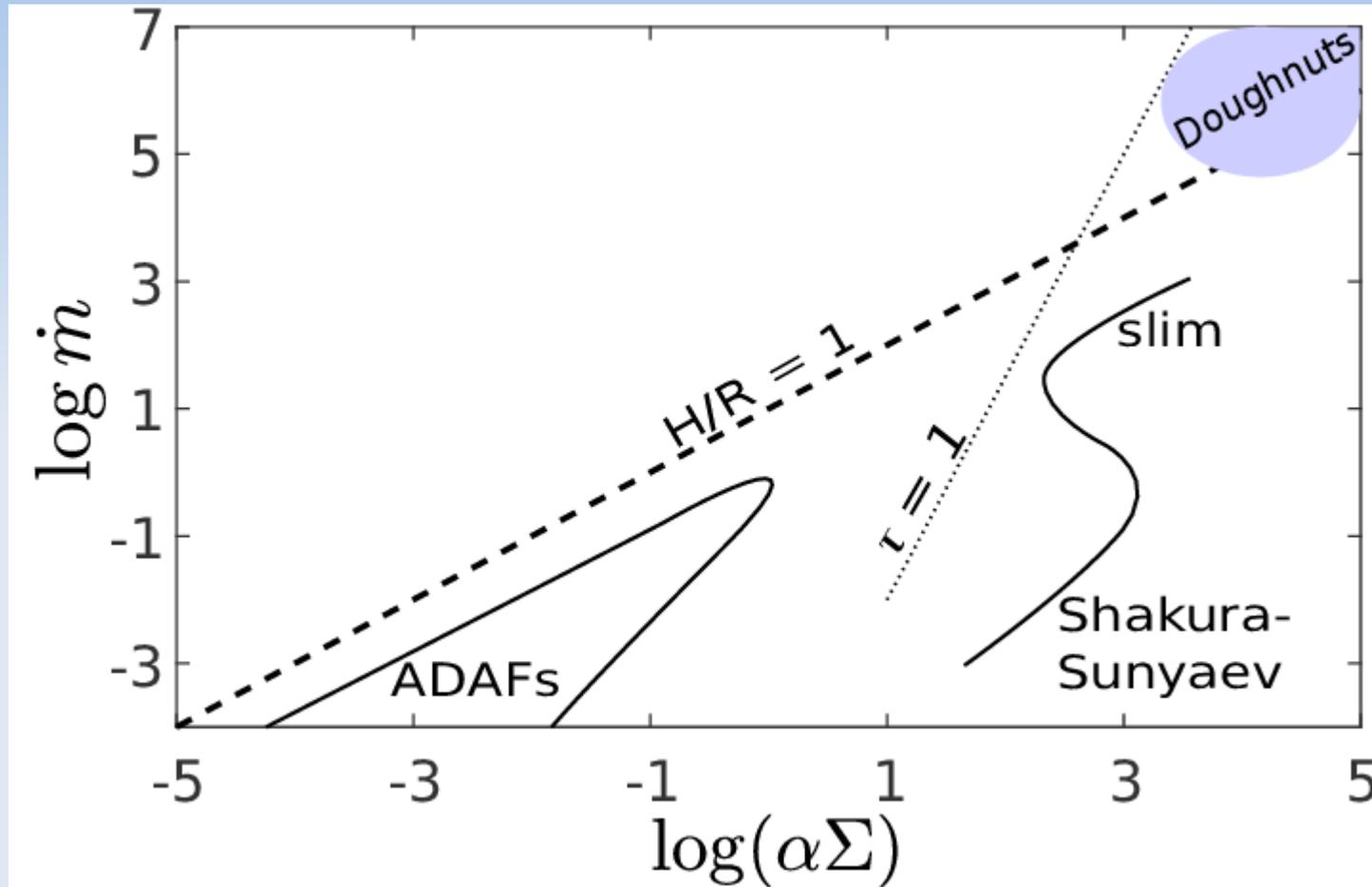


Figure adopted (with some simplifications) from Lasota (2015), the original appeared in Abramowicz et al. (1995).

ADAF branch: developed by Narayan & Yi (1994), cf Yuan (2001)

One-dimensional equations

mass $\frac{d}{dR}(\rho R H v) = 0,$

mom. $v \frac{dv}{dR} - \Omega^2 R = -\Omega_K^2 R - \frac{1}{\rho} \frac{d}{dR}(\rho c_s^2),$

ang. mom. $v \frac{d(\Omega R^2)}{dR} = \frac{1}{\rho R H} \frac{d}{dR} \left(\nu \rho R^3 H \frac{d\Omega}{dR} \right)$

energy
One-T: $\rho v \left(\frac{de}{dR} - \frac{p}{\rho^2} \frac{d\rho}{dR} \right) = \rho \nu R^2 \left(\frac{d\Omega}{dR} \right)^2 - q^-,$

internal E P term viscosity heating radiation cooling

viscous heating collision loss

Two-T: $q^{\text{adv},i} \equiv \rho v \left(\frac{de_i}{dR} - \frac{p_i}{\rho^2} \frac{d\rho}{dR} \right) \equiv \rho v \frac{de_i}{dR} - q^{i,c} = (1 - \delta)q^+ - q^{ie},$

$q^{\text{adv},e} \equiv \rho v \left(\frac{de_e}{dR} - \frac{p_e}{\rho^2} \frac{d\rho}{dR} \right) \equiv \rho v \frac{de_e}{dR} - q^{e,c} = \delta q^+ + q^{ie} - q^-.$

viscous heating collision gain radiation loss

(delta: % of viscous heating that goes to e⁻, unknown, depend on microphysics)



Main features

- Large radial velocity:

$$v_r \sim \frac{\alpha c_s H}{R} \quad c_s = (P/\rho)^{0.5}$$

- Sub-Keplerian rotation: pressure-gradient support

- High temperature: $T \sim \frac{GMm_p}{6kR} \sim \frac{10^{12}}{r}$ (Virial)

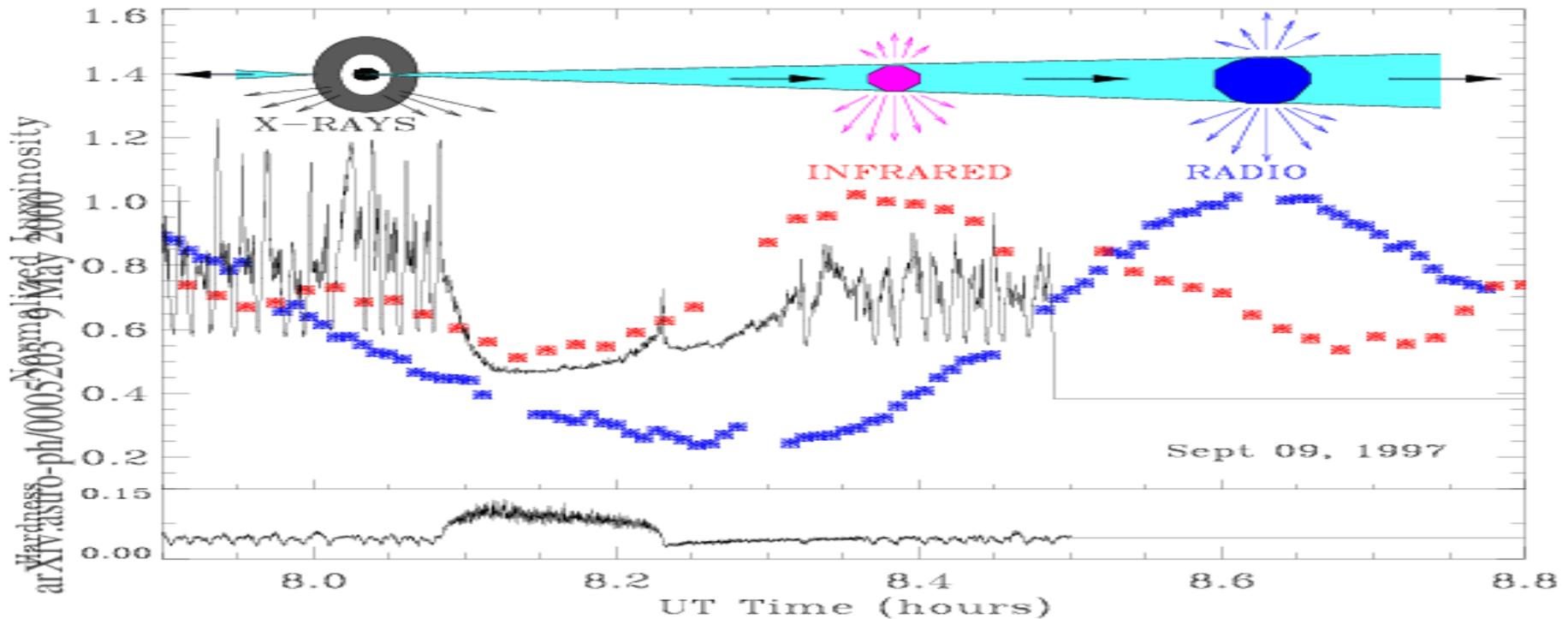
- Geometrically thick: $(H = \frac{c_s}{\Omega_k} \sim R)$ \sim spherical

- Optically thin (because of large radial velocity)

- Two-temperature: $T_i \gg T_e$
 - coupling between ions and electrons not strong enough
 - plasma collective behavior also too weak



Jet-disk connection



Ejection of radio jets are preceded by enhanced X-ray activity.

The jet presence may be related to the accretion disk activity.

Jet-disk connection

- **Is the accretion disk required to produce jets? Probably yes. Collimation is less certain.**
 - For YSO, XRBs, SSS there is a firm observational evidence for the disks' presence
 - For AGNs, the evidence is sometimes far less obvious (emission lines, dusty tori, sub-parsec molecular disks)
 - For PNe, mainly speculative/theoretical arguments for disk presence (Livio & Pringle 1996)
- **Is the jets/outflows required to remove angular momentum from disk?**
 - For some AGN, the MHD driven jets are suggested (observed rotation in winds; but also precession possible)
 - For CVs, the angular momentum is transported in the disk and outer radius expands, wind is not required

What may help driving the jet

- **Central object fast rotation (near break up)**
- **Compactness of the star: determines the escape velocity and so the jet speed**
- High luminosity: radiation pressure driven winds (Proga & Drew 1997)
- Funnel (ADAF flow in AGNs or slow winds in Planetary Nebulae) may provide a collimation
- Boundary layer of NS (?)

Next week

- More about jets physics in general
- How to launch jet from black hole horizon. Blandford Znajek and Penrose processes.
- Morphology of large scale jets from AGN
- Blazars, QSOs, Seyferts: unification scheme

Further reading: R. Fender, "Jets from X-ray binaries", (review), 2006, in: Compact Stellar X-ray sources [arXiv:astro-ph/0303339](https://arxiv.org/abs/astro-ph/0303339)